### Photonic time crystals: A materials 1 perspective 2

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11 Abstract: Recent advances in ultrafast, large-modulation photonic materials have opened the 12 door to many new areas of research. One specific example is the exciting prospect of photonic 13 time crystals. In this perspective, we outline the most recent material advances that are 14 promising candidates for photonic time crystals. We discuss their merit in terms of modulation 15 speed and depth. We also investigate the challenges yet to be faced and provide our estimation 16 on possible roads to success.

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### 19 1. Introduction

20 The realization of Photonic-Time-Crystals (PTCs) holds great potential to unlock new 21 physics and phenomena applicable to non-resonant light amplification, tunable lasing, and a 22 plethora of yet unexplored quantum light-matter interactions [1-10]. In order to realize 23 photonic time-crystals operating in the optical frequency range, several challenges need to be 24 addressed [11]. Most importantly, to achieve a detectable time-reflected signal and/or wide 25 momentum bandgaps, the refractive index of the optical medium must be modulated by an 26 appreciable amount (in the range 0.1-1) all while occurring in a time scale of a single optical 27 cycle of the wave propagating in the medium. Such sharp transitions of permittivity are 28 challenging to realize experimentally. Conventional nonlinear optics, which employs virtual 29 electronic transitions in solids, albeit instantaneous – is orders of magnitude too weak. Other 30 nonlinear optics mechanisms that can yield large index changes typically require transport of some kind, hence are orders of magnitude too slow (photorefractive materials, liquid crystals, 31 32 thermal nonlinearities, etc.). In this perspective, we will first explore possible material 33 platforms for both large and ultrafast index changes. We will then discuss the most viable 34 options for overcoming challenges in both aspects.

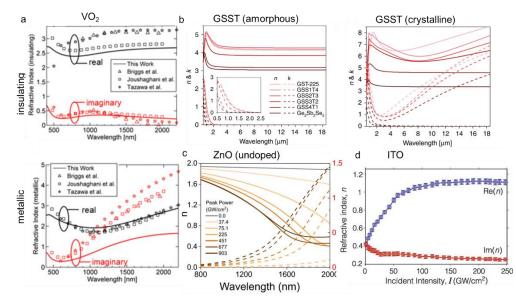
#### 35 2. Materials and mechanisms to trigger large refractive index changes

#### 36 2.1 Large refractive index changes

37 The demanding requirements of achieving an appreciable change in the refractive index 38 within ultra-fast time scales call for extended search for unconventional photonic materials. For 39 different material platforms, various mechanisms have been studied that enable significantly 40 large refractive index change, which at least promises to meet the first criteria for PTC 41 implementation.

42 One such material class is phase change materials that can undergo large refractive index 43 changes as they transition from an amorphous to a crystalline state, and vice versa [12,13]. 44 Vanadium oxide has been shown to have large modulations in kilohertz frequencies under 45 thermal stimulation, with cycle times reaching nanoseconds under an optical impulse [14]. 46 GeSeTe, another prototypical phase change material, has been shown to undergo large 47 relative permittivity changes from 30 in the crystalline phase to 11 in the amorphous 48 phase [15]. Because this transition occurs from thermal changes, the minimum time scale is usually on the order of picoseconds. However, when the optical excitation is below what is 49 50 required for a full phase change, sub-picosecond modulation in the NIR has been observed in 51 both the crystalline [15,16] and amorphous phase [15]. A similar response time of 250 52 femtoseconds was observed in the optical range by P. Martinez et al [17]. In GeSeTe, even if 53 a full phase change is not triggered, the relaxation time is still on the picosecond time scale. 54 Ge–Sb–Se–Te (GSST), another emerging phase change material, shows large changes ( $\Delta n =$ 55 0.5 to 3) (Fig. 1b) in the mid-infrared frequency range, but the nature of the transition is still 56 thermal, with transition times of 30-100 nanoseconds [18].

57 Another promising material class is transparent conducting oxides (TCOs) where the refractive 58 index can be modulated both electrically [19] and optically [20]. Under electrical modulation, 59 an applied voltage accumulates free carriers near the surface of an electrical gate in a TCO, 60 locally changing the refractive index via Drude dispersion [21]. Unity order refractive index 61 changes have been observed in conducting oxides by electrical control [22], with potential 62 speeds reaching the picosecond scale so far [23]. However, despite the large magnitude index 63 changes, the relaxation times with most thermally or electrically activated changes are too slow 64 to implement photonic time crystals. For the same reason, mechanical (MEMS-65 based) [24,25] and structural modulation [18,19] of the effective permittivity are unsuitable for 66 PTC implementation.



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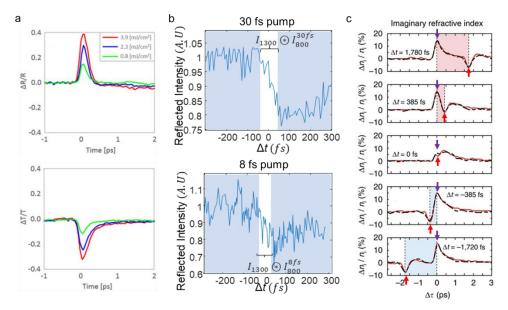
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Figure 1 (a) Vanadium dioxide shows large refractive index changes as it transitions from an insulating to a metallic phase under thermal stimulation [14] (b) GSST can transition from an amorphous to a crystalline stage under an optical pulse, showing unity order changes in the mid infrared wavelength regime. (c) Transparent conducting oxides like zinc oxide have large laser damage threshold, and can tolerate large interband optical pump intensities, enabling large index changes [26](d) Operating near the epsilon-near-zero regime of transparent conducting oxides such as ITO enables large index changes at low pump intensities [20].

75 2.2 Ultrafast index change

76 The second criteria that needs to be met for PTC implementation is the periodic modulation 77 of refractive index, which is more difficult to attain at the required timescales. Since the 78 refractive index change must occur in a time duration comparable to the optical pulse, 79 modulations based on thermal phase-change, which are generally in the nanosecond regime, 80 and electrical changes which are in the picosecond range, are not viable for PTC 81 implementation in the optical frequency domain. Longer wavelengths relax the constraints of 82 the ultrashort cycle time. It is, therefore, no surprise that the first observation of the momentum 83 gap in a time-modulated medium was done at radio frequencies [27] and the first PTC-84 dominated effects such as the exponential growth of EM waves in a temporally varying 85 metasurface, have been demonstrated in the microwave regime [28]. These experiments are 86 also promising signs that for proof-of-concept devices, plasmas or RF circuits can be viable 87 media [29]. Faster nonlinearities in materials such as lithium niobate and barium titanate are 88 another option, but the resistive-capacitive delays of circuits for electrooptic modulation need 89 to be circumvented. Emerging materials such as Gallium phosphide have also been shown to 90 have sub-30 fs optical response times, with large reflectance modulation observed for patterned 91 substrates [30,31], and warrant further exploration in the time-varying field.

92 So far, all-optical excitation seems to be the most promising route to realize PTCs. One 93 potential route for the implementation of PTCs is the optical excitation of free carriers in TCOs. 94 In TCOs, significant changes in the dielectric permittivity can occur much faster than any 95 thermal, thermoelastic, or electrooptic modulation. There are two major pathways to achieve 96 all-optical modulation utilizing the free carriers in TCOs. Firstly, under illumination by an 97 interband pump pulse, electrons jump from the valence band to the conduction band of TCOs, 98 increasing the free carrier concentration and decreasing the permittivity. The material, in this 99 case, can exhibit a large decrease in the refractive index (0.1 in aluminum doped zinc oxide, 100 AZO [32], 1.6 in undoped ZnO [26]). Since the excitation is optical, through absorption of 101 photons, the transition is instantaneous (follows the pump pulse), which can be as fast as a few 102 femtoseconds, and the density of electrons accumulated at the conduction band follows the 103 energy deposited in the medium. The relaxation of electrons back to the valence band, on the 104 other hand, is slower: as the material relaxes through trap-assisted recombination over sub-105 picosecond to few picosecond timescales. However, since the relaxation happens through 106 defect assisted Shockley-Read Hall mechanisms, it can, in principle, be sped up by increasing 107 the defect density in the crystals [33]. In fact, sub-picosecond optical relaxation has been 108 observed in aluminum doped-zinc oxide films operating near epsilon-near-zero (ENZ) (Fig 2a). 109 Altogether, the excitation speed of free carriers in these materials is instantaneous 110 (photoabsorption time), which is limited only by the temporal pulse width of the excitation 111 laser. This allows for ultra-sharp changes in the refractive index upon light absorption (Fig. 2b), which can result in large frequency shifts [34,35], with the potential to achieve time-112 113 reflection [36] under a fast-enough change. In such interband processes the bottleneck is the 114 relaxation speed. However, femtosecond-speed relaxations with intraband pumping are possible, where a sub-bandgap pump energizes conduction electrons, and changing the free 115 116 carrier concentration which leads through an index change. As the carriers relax through 117 thermalization to their original positions in the conduction band, the material returns to the 118 original refractive index. This change can occur in a short time spanning a few hundred 119 femtoseconds [20,37,38]. The use of TCOs has been demonstrated in reflectance and 120 transmittance modulation, as well as in optical polarization switches [13,20,21], making them 121 a promising platform for PTC realization. Furthermore, the large optical nonlinearities near 122 their epsilon-near-zero regimes result in lower power requirements for the same refractive 123 index changes, relaxing the power constraints [20,37].



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Figure 2 (a) Sub picosecond relaxation has been observed in aluminum-doped zinc oxide (AZO), with interband excitation [18] (b) The index change in free-carrier induced modulation is determined by the photoabsorption time (which is instantaneous), and is generally limited by the temporal width of the optical pump. As seen here, using an 8 fs pump instead of a 30 fs pump results in a much sharper optical transition [36] (c) The interband and the intraband nonlinearities can be added. An interband pump (blue arrow) increases the absorption of AZO, and an intraband pump (red arrow) decreases it. It can be seen that if both pumps hit the sample at the same time, the nonlinearities can cancel each other. [39]

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## 134 **3.** Overcoming the limits of fixed relaxation time

135 While optical excitation of materials can be controlled by controlling the pump temporal 136 width, the relaxation mostly depends on electron recombination and lattice cooling, which are 137 intrinsic properties of the materials. Speeding up the relaxation time of materials requires 138 continued exploration of materials with fast enough relaxation time. One viable method of 139 achieving ultrafast changes (positive and negative) with sub-relaxation cycle times, 140 circumventing the intrinsic relaxation time, could be to use two-color excitation. Addition of 141 nonlinearities via concurrent, time-shifted pumping with both intraband and interband pumps 142 has been showed in AZO [39]. This is achieved by using an interband pump to generate photocarriers, decreasing the refractive index, and instantly energizing the carriers to modify 143 144 the index. This might provide a method to obtain a sharper temporal contrast in the permittivity, 145 without relying on the material relaxation time. Another route to achieve ultrafast relaxation 146 time, also with two-color excitations, could be to create population inversion within the 147 conduction band for the lower frequency pump. In this situation, the lower frequency pump 148 could force stimulated emission, which would reduce the density of energetic electrons 149 instantaneously. It is not at all clear that this idea is physical, but it is worth trying.

## 150 4. Overcoming the power constraints

Finally, large and sharp optically induced changes in the refractive index call for millijoules of energy pumped into optical systems in femtosecond cycles, resulting terawatts of peak power [9]. This calls for an in-depth study of materials with large damage thresholds that can simultaneously undergo large index changes. Unity order changes have been observed in doped and undoped oxides with powers 10x lower than the damage threshold of the materials [26,32]. Refractory metals like TiN have been demonstrated to withstand powers over 1 TW [40], for high harmonic generation experiments, and show large reflectance modulation near their 158 epsilon-near-zero regimes [41]. Hybrid systems comprising refractory metals and nitrides may 159 serve as another method of reducing the power constraints. Plasmonic antennae strongly 160 coupled with epsilon-near-zero cavities have been shown to greatly boost nonlinearities due to 161 the enhanced light-matter interaction and slow light effects. Such antennae have been shown to 162 increase the effective optical nonlinearity [42], and even lower the power requirements for time 163 refraction [8,32] and photon acceleration [43] in ENZ-nanoantenna systems, making them a 164 promising route to alleviate the high-power requirements. In such system, robust epsilon-near-165 zero substrates can be used to increase the nonlinear interaction between the incident wave and 166 antennae made of refractory metals, boosting the time-reflected signal.

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# 168 **5. Conclusion**

169 While the ultrafast and large modulation of the material's optical response remains a 170 challenge for the photonic time-crystal implementation, the range of materials that allow for a 171 fast and substantial change of their refractive index is rapidly growing. At present times, the 172 most promising class of materials that offers large optical response changes, fast modulation, 173 and large damage threshold is transparent conducting oxides operating near their epsilon near 174 zero regime. While TCOs promise to be viable candidates to implement PTCs, emerging 175 materials such as monolayer MoSe<sub>2</sub> [44], and transdimensional films [45], also show promise 176 in terms of speed and modulation depth. Techniques to reduce the modulation power 177 requirements while sustaining PTC effects call for both rigorous theoretical study of different 178 schemes of refractive index modulation, with a continued quest for finding better, faster, and 179 more laser tolerant materials.

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# 186 Disclosures:

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- **187** The authors declare no conflicts of interest.
- 189 **Data availability.** No data were generated or analyzed in the presented 190 research.

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